

Universidade do Minho Centro de Matemática

Additive functionals of exclusion processes

PATRÍCIA GONÇALVES Centro de Matemática da Universidade do Minho

Braga (Portugal) **Interacting Particle Systems and Related Topics**

Florence, Italy

August 2012 (joint work with Milton Jara - IMPA)





1. Abstract

I consider exclusion processes evolving on \mathbb{Z} and starting from the invariant state. The goal consists in establishing scaling limits of $\Gamma_t(f) := \int_0^t f(\eta_s) ds$ for local functions f. I present a method, recently introduced in [3], from which we derive a local Boltzmann-Gibbs Principle for a class of exclusion processes. For the occupation time of the origin (i.e. for $f(\eta) := \eta(0)$), this principle says that $\Gamma_t(f)$ is very well approximated to the density of particles. As a consequence, the scaling limits of $\Gamma_t(f)$ follow from the scaling limits of the density of particles. As examples I present the mean-zero exclusion, the symmetric simple exclusion and the weakly asymmetric simple exclusion. For the latter under a strong asymmetry, we establish the limit of $\Gamma_t(f)$ in terms of the solution of the KPZ equation.

5. Strategy of the proof

Theorem 3: Local Boltzmann-Gibbs principle. Le $f: \Omega \to \mathbb{R}$ be a local function, such that $supp(f) \subseteq \{1, ..., k\}$ and $\varphi_f(\rho) = 0$. There exists $c = c(f, \rho)$ such that

i) if $\varphi'_f(\rho) \neq 0$, then for any $t \geq 0$ and any $\ell \geq k$:

 $\mathbb{E}_{\rho}\Big[\Big(\int_{0}^{t} \Big\{f(\eta_{s}) - \varphi_{f}'(\rho)\big(\eta_{s}^{\ell} - \rho\big)\Big\}ds\Big)^{2}\Big] \leq c\Big(t\ell + \frac{t^{2}}{\ell^{2}}\Big),$

ii) if $\varphi'_f(\rho) = 0$, then for any $t \ge 0$ and any $\ell \ge k$:

The measures $\{\nu_{\rho}; \rho \in [0, 1]\}$ are **invariant**, but they are **not necessarily reversible** (that is true if and only if $p(\cdot)$ is symmetric). The mean-zero, non-symmetric exclusion process is a **diffusive** and **non-reversible system**.

We can define the density fluctuation field $\{\mathcal{Y}_t^n; t \in [0, T]\}$ as in (2) and we have that:

Proposition 2: The process $\{\mathcal{Y}_t^n; t \in [0,T]\}$ converges in distribution with respect to the Skorohod topology of $\mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$ to the stationary solution of the Ornstein-Uhlenbeck equation

 $d\mathcal{Y}_t = D(\rho)\Delta\mathcal{Y}_t dt + \sqrt{2D(\rho)\rho(1-\rho)(\rho)}\nabla d\mathcal{B}_t,$

where $D(\rho)$ is the diffusion coefficient.

The results presented above allow us to get the scaling limits of additive functionals as in Theorem 2.

2-Symmetric simple exclusion

2. Lattice gas dynamics

• η_t is a Markov process with space state $\Omega := \{0, 1\}^{\mathbb{Z}}$. • for $x \in \mathbb{Z}$, $\eta(x) = 1$ if the site x is occupied, otherwise $\eta(x) = 0$. • Let $r: \Omega \to \mathbb{R}$ be a local function that satisfies: i) There exists $\varepsilon_0 > 0$ such that $\varepsilon_0 < r(\eta) < \varepsilon_0^{-1}$ for any $\eta \in \{0, 1\}^{\mathbb{Z}}$. (Ellipticity) ii) For any η , ξ such that $\eta(x) = \xi(x)$ for $x \neq 0, 1$, then $r(\eta) = r(\xi)$. (**Reversibility**) • The generator is given on local functions $f: \Omega \to \mathbb{R}$ by:

 $\mathcal{L}f(\eta) = \sum_{x \in \mathbb{Z}} r(\tau_x \eta) (f(\eta^{x,x+1}) - f(\eta))$

where

 $\eta^{x,y}(z) = \begin{cases} \eta(y), & z = x\\ \eta(x), & z = y\\ \eta(z), & z \neq x, y. \end{cases}$

(1)

(2)

(5)

Invariant states

Let ν_{ρ} be the Bernoulli product measures of constant parameter $\rho \in [0, 1]$. Under this measure the occupation variables $(\eta(x))_x$ are independent and

 $\nu_{\rho}(\eta:\eta(x)=1)=\rho.$

We denote by \mathbb{E}_{ρ} the expectation with respect to \mathbb{P}_{ρ} - the distribution of $\{\eta_t : t \geq 0\}$ in the space $\mathcal{D}([0,\infty],\Omega)$ starting from ν_{ρ} .

3. Scaling Limits of the density of particles

 $\mathbb{E}_{\rho}\Big[\Big(\int_{0}^{t} \Big\{f(\eta_{s}) - \frac{\varphi_{f}''(\rho)}{2}\Big(\big(\eta_{s}^{\ell} - \rho\big)^{2} - \frac{\rho(1-\rho)}{\ell}\Big)\Big\}ds\Big)^{2}\Big] \le c\Big(t(\log\ell)^{2} + \frac{t^{2}}{\ell^{3}}\Big)$ where $\eta^{\ell} := \frac{1}{\ell} \sum_{x=1}^{\ell} \eta(x)$.

The proof of last result is divided into four steps that rely on the Kipnis-Varadhan inequality (see [5]) and the spectral gap inequality (see [6]): **1.** Firstly, we compare the additive functional associated to f with the additive functional associated to $\psi_f(\ell) := E_\rho[f|\sum_{x=1}^\ell \eta(x)]$, using the:

Lemma 1: One-block estimate. Let $f : \Omega \to \mathbb{R}$ be a local function such that $\varphi_f(\rho) = 0$. Then, there exists $c = c(f, \rho)$ such that for any $\ell \geq k$ and any $t \ge 0$: $\mathbb{E}_{\rho}\left[\left(\int_{0}^{\iota} \{f(\eta_{s}) - \psi_{f}(\ell;\eta_{s})\}ds\right)^{2}\right] \leq ct\ell^{2} \operatorname{Var}(f;\nu_{\rho}).$

2. Secondly, we compare the additive functional associated to $\psi_f(\ell)$ with the additive functional associated to $\psi_f(2\ell)$, using the:

Lemma 2: Renormalization step. Let $f : \Omega \to \mathbb{R}$ be a local function such that $\varphi_f(\rho) = 0$. There exists $c = c(f, \rho)$ such that for any $\ell \geq k$ and any $t \ge 0$:

$$\mathbb{E}_{\rho}\Big[\Big(\int_{0}^{t} \{\psi_{f}(\ell;\eta_{s}) - \psi_{f}(2\ell;\eta_{s})\}ds\Big)^{2}\Big] \leq \begin{cases} ct\ell, if \ \varphi_{f}'(\rho) \neq 0, \\ ct, if \ \varphi_{f}'(\rho) = 0. \end{cases}$$

3. Thirdly, we compare the additive functional associated to $\psi_f(k)$ with the additive functional associated to $\psi_f(2^m k)$, using the renormalization step m times.

Lemma 3: Two-blocks estimate. Let $f : \Omega \to \mathbb{R}$ be a local function such that $\varphi_f(\rho) = 0$. Then, there exists $c = c(f, \rho)$ such that for any $\ell \geq k$ and any $t \ge 0$:

Let p be such that p(1) = p(-1) = 1/2. The measures $\{\nu_{\rho}; \rho \in [0, 1]\}$ are **invariant**.

Proposition 3: The process $\{\mathcal{Y}_t^n; t \in [0, T]\}$ converges in distribution with respect to the Skorohod topology of $\mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$ to the stationary solution of the Ornstein-Uhlenbeck equation

$$d\mathcal{Y}_t = \frac{1}{2}\Delta \mathcal{Y}_t dt + \sqrt{\rho(1-\rho)}\nabla d\mathcal{B}_t.$$

The results presented above allow us to get the scaling limits of additive functionals as in Theorem 2.

3-The weakly asymmetric simple exclusion

Let $p_n(1) = \frac{1}{2} + \frac{a_n}{2}$, $p_n(-1) = \frac{1}{2} - \frac{a_n}{2}$ and $p_n(z) = 0$ if $z \neq -1, 1$. The measures $\{\nu_\rho; \rho \in [0, 1]\}$ are **invariant**. If $a_n := \frac{1}{n}$, then we have that

Proposition 4: The process $\{\mathcal{Y}_t^n; t \in [0,T]\}$ converges in distribution with respect to the Skorohod topology of $\mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$ to the stationary solution of the Ornstein-Uhlenbeck equation

$$d\mathcal{Y}_t = \frac{1}{2}\Delta \mathcal{Y}_t dt + (1-2\rho)\nabla \mathcal{Y}_t dt + \sqrt{\rho(1-\rho)}\nabla d\mathcal{B}_t.$$

In this case the Ornstein-Uhlenbeck process has a drift, nevertheless one can get the same result as in Theorem 2.

4-The weakly asymmetric simple exclusion, the KPZ scaling

If $\rho = 1/2$, then the limit field is the same as in the symmetric simple exclusion (so a weak asymmetry does not have influence!), see [4]. In this case the "correct" strength asymme**try** is $a_n = 1/\sqrt{n}$. In this case we have

Proposition 5: The process $\{\mathcal{Y}_t^n; t \in [0, T]\}$ converges in distribution with respect to the Skorohod topology of $\mathcal{D}([0,T]; \mathcal{S}'(\mathbb{R}))$ to the stationary solution of the stochastic Burgers equation $d\mathcal{Y}_t = \frac{1}{2}\Delta \mathcal{Y}_t dt + \left(\nabla \mathcal{Y}_t\right)^2 dt + \sqrt{\rho(1-\rho)}\nabla d\mathcal{B}_t.$

• Hydrodynamic Limit For each configuration η we denote by $\pi^n(\eta; du)$ the empirical measure:

 $\pi^n(\eta; du) = \frac{1}{n} \sum_{x \in \mathbb{Z}} \eta(x) \delta_{x/n}$

and $\pi_t^n(\eta, du) := \pi^n(\eta_t, du).$ Under a diffusive scaling of time tn^2 , the hydrodynamic limit (LLN) was obtained by [2].

• Equilibrium fluctuations:

The density fluctuation field acts on functions $G \in \mathcal{S}(\mathbb{R})$ as

 $\mathcal{Y}_t^n(G) := \frac{1}{\sqrt{n}} \sum_{\pi} G\left(\frac{x}{n}\right) \{\eta_{tn^2}(x) - \rho\}.$

It was proved by [1] that $\{\mathcal{Y}_t^n; t \in [0,T]\}$ converges in distribution with respect to the Skorohod topology of $\mathcal{D}([0,T], \mathcal{S}'(\mathbb{R}))$ to the stationary solution of the Ornstein-Uhlenbeck equation (3)

 $d\mathcal{Y}_t = D(\rho)\Delta\mathcal{Y}_t dt + \sqrt{2D(\rho)\rho(1-\rho)}\nabla d\mathcal{B}_t,$

where \mathcal{B}_t is a $\mathcal{S}'(\mathbb{R})$ -valued Brownian motion and $D(\rho)$ is the diffusion coefficient.

This means that the trajectories of the limit field \mathcal{Y}_t are in $\mathcal{C}([0,T], \mathcal{S}'(\mathbb{R}))$ and that \mathcal{Y}_0 is a white noise of variance $\rho(1-\rho)$ - namely if for any $G \in \mathcal{S}(\mathbb{R})$, the real-valued random variable $\mathcal{Y}_0(G)$ has a normal distribution of mean zero and variance $\rho(1-\rho) \|G\|^2$.

Now, fix a stationary solution $\{\mathcal{Y}_t; t \in [0, T]\}$ of (3). For $x \in \mathbb{R}$, let $i_{\varepsilon}(x) : y \mapsto \varepsilon^{-1} \mathbf{1}_{(0,1]}((y-x)\varepsilon^{-1})$.

Theorem 1: For each $\varepsilon \in (0, 1)$, let $\{\mathcal{Z}_t^{\varepsilon}; t \in [0, T]\}$ be defined as $\mathcal{Z}_t^{arepsilon} = \int_0^t \mathcal{Y}_s(i_arepsilon) ds.$

Then, the process $\{\mathcal{Z}_t^{\varepsilon}; t \in [0,T]\}$ converges in distribution with respect to the

 $\mathbb{E}_{\rho}\left[\left(\int_{0}^{t}\psi_{f}(k;\eta_{s})-\psi_{f}(\ell;\eta_{s})ds\right)^{2}\right] \leq \begin{cases} ct\ell, if \ \varphi_{f}'(\rho) \neq 0, \\ ct(\log \ell)^{2}, if \ \varphi_{f}'(\rho) = 0. \end{cases}$

4. Finally, we replace $\psi_f(\ell)$ by the corresponding function of η^{ℓ} using the following:

Proposition 1: Equivalence of Ensembles. Let $f : \Omega \rightarrow I$ \mathbb{R} be a local function. Then there exists a constant $c = c(f, \rho)$ such that for any $\ell \geq k$:

 $\int \left(\psi_f(\ell,\eta) - \varphi_f'(\rho) \left(\eta^\ell - \rho\right) - \frac{\varphi_f''(\rho)}{2} \left(\left(\eta_s^\ell - \rho\right)^2 - \frac{\rho(1-\rho)}{\ell}\right)\right)^2 d\nu_\rho \le \frac{c}{\ell^3}.$

6. Examples

1-Mean-Zero Exclusion

Initially one can have a certain number of particles, as for example:

After an exponential time of parameter 1, a particle at x jumps to x + y with probability p(y), but to respect the exclusion rule, particles can only jump to empty sites. So, this jump is allowed:

Theorem 4: Let $\{\mathcal{Y}_t; t \in [0,T]\}$ the stationary solution of the stochastic Burgers equation above. For $\varepsilon > 0$, let $\mathcal{Z}_t^{\varepsilon} = \int_0^t \mathcal{Y}_s(i_{\varepsilon}) ds$. Then there exists $\{\mathcal{Z}_t; t \in [0,T]\}$ such that $\{\mathcal{Z}_t^{\varepsilon}; t \in [0,T]\}$ converges in distribution with respect to the uniform topology of $\mathcal{C}([0,T];\mathbb{R})$, as $\varepsilon \to 0$, to $\{\mathcal{Z}_t; t \in [0,T]\}$.

And as a consequence we have that

Theorem 5: Let $f: \Omega \to \mathbb{R}$ be a local function such that $\varphi_f(1/2) = 0$. Then, $\{\Gamma_t^n(f); t \in [0,T]\}$ as defined in (5) converges in distribution with respect to the uniform topology of $\mathcal{C}([0,T];\mathbb{R})$ to $\{\varphi'_f(1/2)\mathcal{Z}_t; t \in [0,T]\}$, where \mathcal{Z}_t is the same as in Theorem 4.

5-Symmetric simple exclusion/Asymmetric simple exclusion

Let f be a local function. Then

• if $\varphi_f(\rho) = 0, \ \varphi'_f(\rho) \neq 0$, then $\operatorname{Var}(\Gamma_t(f); \nu_{\rho}) \leq t^{3/2}(t^{4/3}, \rho = 1/2; t, \rho \neq 1/2)$,

With the results presented above we can get the correct upper bound in the symmetric case. Our method does not give the correct upper bound in the asymmetric case. • if $\varphi_f(\rho) = \varphi'_f(\rho) = 0, \ \varphi''_f(\rho) \neq 0$, then $\operatorname{Var}(\Gamma_t(f); \nu_\rho) \leq t \log(t)(t)$,

With the results presented above we can get the upper bound $t(log(t))^2$ in the symmetric case. Our method does not give the correct upper bound in the asymmetric case.

• if $\varphi_f(\rho) = \varphi'_f(\rho) = \varphi''_f(\rho) = 0, \ \varphi'''_f(\rho) \neq 0$, then $\operatorname{Var}(\Gamma_t(f); \nu_\rho) \leq t$.

uniform topology of $\mathcal{C}([0,T],\mathbb{R})$, as $\varepsilon \to 0$, to a fractional Brownian motion $\{\mathcal{Z}_t; t \in [0, T]\}$ of Hurst exponent H = 3/4.

4. Additive functionals

Our goal consists in obtaining functional limit theorems for observables of the process $\{\eta_t; t \geq 0\}$. The occupation time of a site $x \in \mathbb{Z}$ is defined as the integral $\int_0^t \eta_s(x) ds$. More generally, for $f: \Omega \to \mathbb{R}$ a local function, if for $\beta \in [0, 1], \varphi_f(\beta) = \int f d\nu_\beta$, then:

Theorem 2: The process $\{\Gamma_t^n(f); t \in [0, T]\}$ defined as

$${n \choose t} f = rac{1}{n^{3/2}} \int_0^{tn^2} \left(f(\eta_s) - \varphi_f(\rho) \right) ds$$

converges in distribution with respect to the uniform topology of $\mathcal{C}([0,T],\mathbb{R})$ to $\{\varphi'_f(\rho)\mathcal{Z}_t; t \in [0,T]\}, \text{ where } \{\mathcal{Z}_t; t \in [0,T]\} \text{ is the same as above.}$

But, this jump is forbidden:



The microscopic dynamics:

- Let $p : \mathbb{Z} \setminus \{0\} \to [0, 1]$ be a probability measure, such that:
- 1. $p(\cdot)$ has finite range, that is, there exists M > 0 such that p(z) = 0 whenever |z| > M; 2. $p(\cdot)$ is irreducible, i.e. $\mathbb{Z} = \operatorname{span}\{z \in \mathbb{Z}; p(z) > 0\};$ 3. $p(\cdot)$ has mean zero: $\sum_{z \in \mathbb{Z}} zp(z) = 0$.

Example: p(1) = 2/3, p(-2) = 1/3 and p(z) = 0 if $z \neq -2, 1$.

We define the Markov process $\{\eta_t^{\text{ex}}; t \geq 0\}$ whose **generator** acts over local functions $f: \Omega \to \mathbb{R}$ as

 $\mathcal{L}_{\mathrm{ex}}f(\eta) = \sum_{x,y\in\mathbb{Z}} p(y)\eta(x)(1-\eta(x+y))(f(\eta^{x,x+y})-f(\eta)),$

with $p(\cdot)$ as above and $\eta^{x,x+y}$ as in (1).

With the results presented	d above	we can	get the	correct	upper	bound i	in both	cases,	see
also $[7]$.									

7. References

- [1] C. Chang. Equilibrium fluctuations of gradient reversible particle systems. *Probability* Theory and Related Fields, 1994.
- [2] T. Funaki, K. Uchiyama, and H. Yau. Hydrodynamic limit for lattice gas reversible under Bernoulli measures. In Nonlinear stochastic PDEs, Springer, 1996.
- [3] Gonçalves, P. and Jara, M. (2011): Universality of KPZ equation, available online at arXiv:1003.4478.
- [4] Gonçalves, P and Jara, M. (2011): Crossover to the KPZ equation, Annales Henri Poincaré, Volume 13, Number 4, 813-826.
- [5] C. Kipnis and S. Varadhan, Central limit theorem for additive functionals of reversible Markov processes and applications to simple exclusions. Communications in Mathematical Physics 1986.
- [6] J. Quastel. Diffusion of color in the simple exclusion process, *Communications on Pure* and Applied Mathematics, 1992.
- [7] S. Sethuraman, L. Xu. A central limit theorem for reversible exclusion and zero-range particle systems. Annals of Probability, 1996.